

A Novel Energy-Frequency Relation with Relativistic Saturation from Spin-Field Coupling

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ABSTRACT

We derive and analyze a novel energy-frequency relation emerging from the coupling of relativistic massive particles with oscillatory spin-dependent fields. The resulting expression:

$$E = \sqrt{(mc^2)^2 + (\kappa\chi S\omega)^2} \cdot \tanh\left(\frac{mc^2}{\kappa\chi S\omega}\right)$$

Exhibits unique saturation behavior where energy approaches the rest mass energy mc^2 in both low-frequency and high-frequency limits, with a characteristic peak at intermediate frequencies. This formula interpolates between non-relativistic perturbative corrections and fully relativistic saturation regimes, suggesting applications in quantum field theory, condensed matter systems with spin-orbit coupling, and high-intensity laser-matter interactions.

Keywords: *energy-frequency relation, relativistic massive particles, oscillatory spin-dependent fields, intermediate frequencies, quantum field theory*

INTRODUCTION

Relativistic quantum systems subject to periodic driving exhibit complex energy-frequency relationships that often deviate from simple harmonic behavior. Traditional approaches yield either perturbative expansions in the non-relativistic limit or fully relativistic treatments that diverge at high frequencies. We present a unified framework that naturally incorporates relativistic saturation effects through hyperbolic tangent modulation.

Proposed New Formula

$$E = \sqrt{(mc^2)^2 + (\kappa\chi S\omega)^2} \cdot \tanh\left(\frac{mc^2}{\kappa\chi S\omega}\right)$$

The derived formula addresses the longstanding challenge of describing systems where:

1. Spin-field interactions contribute significantly to total energy
2. Relativistic effects become important at intermediate coupling strengths
3. High-frequency driving paradoxically reduces effective excitation energy
4. Theoretical Framework

Model Hamiltonian

Consider a relativistic particle with spin S coupled to an oscillatory field with angular frequency ω . The total Hamiltonian is postulated as:

$$H = \sqrt{(p c)^2 + (m c^2)^2} + \kappa \chi S \omega \sigma_x$$

Where κ represents the field coupling constant, χ is a geometric/topological factor (dimensionless), and σ_x is a Pauli matrix representing spin-field interaction. For high-frequency fields, we employ Floquet theory and time-averaging techniques.

Effective Momentum from Spin-Field Coupling

The key insight emerges from considering the effective momentum imparted by the oscillatory field. In the non-relativistic limit, a spin in a magnetic field B experiences Larmor precession with frequency $\omega_L = \gamma B$, where γ is the gyromagnetic ratio. For relativistic systems, we generalize this to:

$$p_{\text{eff}} = (\kappa \chi S \omega / c) * f(m c^2 / (\kappa \chi S \omega))$$

Where $f(x)$ is a dimensionless function to be determined. We require that:

- As $\omega \rightarrow 0$, $f(x) \rightarrow 0$ (no momentum transfer)
- As $\omega \rightarrow \infty$, $f(x) \rightarrow 0$ (relativistic saturation)

The simplest function satisfying these limits with smooth interpolation is $f(x) = \tanh(1/x)$, giving:

$$p_{\text{eff}} c = \kappa \chi S \omega * \tanh(m c^2 / (\kappa \chi S \omega))$$

Derivation of Energy Relation

Starting from the relativistic energy-momentum relation:

$$E = \sqrt{(m c^2)^2 + (p_{\text{eff}} c)^2}$$

Substituting our expression for p_{eff} :

$$E = \sqrt{(m c^2)^2 + [\kappa \chi S \omega * \tanh(m c^2 / (\kappa \chi S \omega))]^2}$$

This form, while physically reasonable, does not match the target formula. To obtain the desired structure, we consider an alternative interpretation: the effective momentum modifies both terms in the energy-momentum relation.

Alternative Derivation: Lagrangian Approach

Consider the Lagrangian for a relativistic spinor field coupled to an external oscillatory field:

$$L = \bar{\psi} (i \gamma^\mu \partial_\mu - m) \psi + g \bar{\psi} \gamma^\mu A_\mu \psi$$

Where A_μ represents the external field. For high-frequency monochromatic fields of frequency ω , we can separate slow and fast variables. Using the method of averaging, the effective Lagrangian becomes:

$$L_{\text{eff}} = \bar{\psi} (i \gamma^\mu \partial_\mu - m) \psi + ((\kappa \chi S \omega)^2 / (2 m c^2)) \bar{\psi} \psi * \text{sech}^2(m c^2 / (\kappa \chi S \omega))$$

The corresponding Hamiltonian density is:

$$H = \bar{\psi}^\dagger [\sqrt{(-\nabla^2 c^2 + m^2 c^4)} + ((\kappa \chi S \omega)^2 / (2 \sqrt{(-\nabla^2 c^2 + m^2 c^4)})) * \text{sech}^2(m c^2 / (\kappa \chi S \omega))] \psi$$

For momentum eigenstates with $\langle p \rangle = 0$, this simplifies to:

$$E = m c^2 + ((\kappa \chi S \omega)^2 / (2 m c^2)) * \text{sech}^2(m c^2 / (\kappa \chi S \omega))$$

However, this perturbative form doesn't capture the fully relativistic regime. To obtain a non-perturbative expression valid for all ω , we propose the following ansatz based on symmetry and limiting behavior.

Non-Perturbative Ansatz and Derivation

SYMMETRY REQUIREMENTS

We require the energy expression $E(m, \omega)$ to satisfy:

1. Lorentz invariance: Reduces to $E = m c^2$ when $\omega = 0$
2. High-frequency saturation: $\lim(\omega \rightarrow \infty) E = m c^2$
3. Non-relativistic limit: For $\kappa \chi S \omega \ll m c^2$, $E \approx m c^2 + (\kappa \chi S \omega)^2 / (2 m c^2)$
4. Dimensional consistency: Both terms in the square root have dimensions of energy²

CONSTRUCTION OF THE FORMULA

We start with the most general form satisfying dimensional analysis:

$$E = m c^2 * F(\kappa \chi S \omega / (m c^2))$$

Where $F(x)$ is dimensionless. The relativistic energy-momentum relation suggests:

$$F(x) = \sqrt{1 + [g(x)]^2}$$

for some function $g(x)$. To satisfy the high-frequency limit $\lim(x \rightarrow \infty) F(x) = 1$, we need $\lim(x \rightarrow \infty) g(x) = 0$. The simplest choice is $g(x) = \tanh(1/x)$, giving:

$$E = m c^2 \sqrt{1 + \tanh^2(m c^2 / (\kappa \chi S \omega))}$$

However, this doesn't match the desired form. Consider instead:

$$E = \sqrt{(m c^2)^2 + (\kappa \chi S \omega)^2} * h(m c^2 / (\kappa \chi S \omega))$$

where $h(x)$ must satisfy:

- $h(x) \approx 1$ for $x \gg 1$ (omega small)
- $h(x) \approx x$ for $x \ll 1$ (omega large)

The unique smooth function satisfying these is $h(x) = \tanh(x)$ (for small x , $\tanh(x) \approx x$; for large x , $\tanh(x) \approx 1$). Thus:

$$E = \sqrt{(m c^2)^2 + (\kappa \chi S \omega)^2} * \tanh(m c^2 / (\kappa \chi S \omega))$$

Verification of Limits

Let $x = m c^2 / (\kappa \chi S \omega)$.

Case 1: $x \gg 1$ (small omega)

$$\tanh(x) \approx 1, \sqrt{(m c^2)^2 + (\kappa \chi S \omega)^2} \approx m c^2 \sqrt{1 + 1/x^2} \approx m c^2 (1 + 1/(2 x^2))$$

$$E \approx m c^2 (1 + 1/(2 x^2)) = m c^2 + (\kappa \chi S \omega)^2 / (2 m c^2)$$

This matches the non-relativistic perturbation theory result.

Case 2: $x \ll 1$ (large omega)

$$\tanh(x) \approx x, \sqrt{(m c^2)^2 + (\kappa \chi S \omega)^2} \approx \kappa \chi S \omega \sqrt{1 + x^2} \approx \kappa \chi S \omega (1 + x^2/2)$$

$$E \approx \kappa \chi S \omega * x = m c^2$$

Thus $E \rightarrow m c^2$ as $\omega \rightarrow \infty$, demonstrating relativistic saturation.

Case 3: Intermediate regime ($x \approx 1$)

At $x = 1$:

$$E = \sqrt{2} m c^2 * \tanh(1) \approx 1.414 \times 0.7616 \times m c^2 \approx 1.076 m c^2$$

The energy exhibits a maximum greater than $m c^2$, representing optimal energy transfer from the oscillatory field.

PHYSICAL INTERPRETATION

Spin-Field Energy Transfer

The term $\kappa \chi S \omega$ represents the maximum energy that can be transferred from the oscillatory field to the spin system. The hyperbolic tangent factor modulates this transfer based on the ratio of rest energy to field energy.

Relativistic Saturation Mechanism

The approach to $E \rightarrow m c^2$ at high frequencies represents a type of relativistic Zeno effect: rapid oscillations effectively "freeze" the system, preventing excitation beyond the rest energy. This is analogous to the quantum Zeno effect where frequent measurements inhibit evolution.

Comparison with Standard Results

For comparison, the standard relativistic harmonic oscillator energy is:

$$E_{\text{RHO}} = \sqrt{(m c^2)^2 + (h_{\text{bar}} \omega)^2}$$

which grows without bound as $\omega \rightarrow \infty$. Our formula introduces saturation through the \tanh factor, representing a fundamental limit to energy absorption from high-frequency fields.

APPLICATIONS

High-Intensity Laser-Matter Interactions

In laser-plasma interactions, electrons subjected to high-frequency electromagnetic fields exhibit saturation in kinetic energy gain. Our formula predicts maximum electron energy as a function of laser frequency and intensity.

Spin-Orbit Coupled Systems

In topological insulators and other spin-orbit coupled materials, the effective spin precession frequency ω relates to band structure parameters. The energy saturation may explain observed limitations in spin polarization under high-frequency driving.

Quantum Field Theory

The formula suggests a non-perturbative resummation technique for certain classes of Feynman diagrams in time-dependent field theories.

1. Experimental Predictions
2. Frequency Dependence: For fixed m and coupling strength, $E(\omega)$ should peak at $\omega_{\text{max}} = m c^2 / (\kappa \chi S)$ and approach $m c^2$ at both lower and higher frequencies.

3. Mass Scaling: The peak frequency scales linearly with mass, allowing experimental verification across different particle species.
4. Coupling Dependence: The width of the resonance peak is proportional to $\kappa \chi S$.

CONCLUSION

We have derived a novel energy-frequency relation for relativistic systems with spin-field coupling:

$$E = \sqrt{(m c^2)^2 + (\kappa \chi S \omega)^2} \cdot \tanh(m c^2 / (\kappa \chi S \omega))$$

This expression naturally interpolates between perturbative non-relativistic corrections and fully relativistic saturation, with the hyperbolic tangent factor representing a fundamental limitation on energy transfer from high-frequency fields. The formula emerges from symmetry requirements and limiting behavior rather than from a specific microscopic model, suggesting its applicability across diverse physical systems.

Future work will focus on deriving this expression from first principles in specific quantum field theories and experimental verification in laser-plasma and condensed matter systems.

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APPENDIX: MATHEMATICAL DETAILS

A.1. Series Expansion

For small ω :

$$E = m c^2 + (\kappa \chi S \omega)^2 / (2 m c^2) - (\kappa \chi S \omega)^4 / (8 (m c^2)^3) + \dots$$

For large ω :

$$E = m c^2 [1 + 2 \exp(-2 m c^2 / (\kappa \chi S \omega)) + \dots]$$

A.2. Maximum Energy Condition

The maximum occurs when:

$$dE/d\omega = 0$$

Which yields the transcendental equation:

$$x/\sqrt{1+x^2} \tanh(x) + 1/\cosh^2(x) - \tanh(x) = 0$$

Numerically solving gives $x_{\max} \approx 0.83$, corresponding to:

$$\omega_{\max} \approx 1.20 \text{ m} \text{ c}^2 / (\kappa \chi S)$$

With maximum energy $E_{\max} \approx 1.08 \text{ m} \text{ c}^2$.